

Solving the KPZ equation

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Introduction

Object of study: KPZ equation of surface growth:

$$\partial_t h = \partial_x^2 h + \lambda (\partial_x h)^2 + \xi - \infty$$

with either $x \in \mathbf{R}$ or $x \in S^1$, and ξ is **space-time** white noise.

1. Universal model for interface fluctuations. (Shown rigorously only for SOS model, see Bertini-Giacomin 1997.)
2. Free energy for polymer models.
3. Scaling limit of time-dependent parabolic Anderson model.
4. Universal object describing crossover from Edwards-Wilkinson to KPZ.

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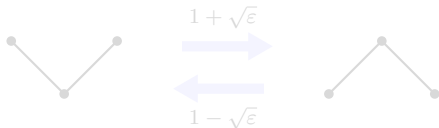
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The SOS model

Simplest possible model of surface growth. Surface modelled by graph with slope ± 1 :



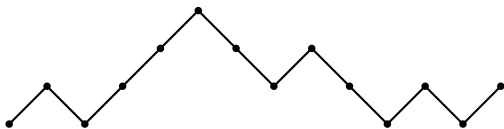
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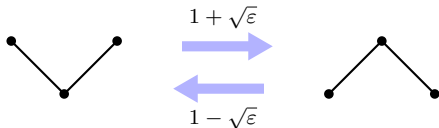
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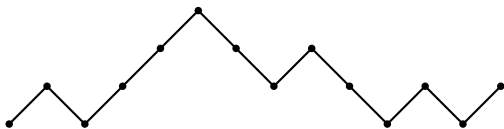
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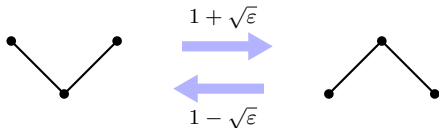
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Cole-Hopf solution

Trick introduced by Cole and Hopf in the 50's. Write

$$h = \lambda^{-1} \log Z ,$$

then Z solves

$$\partial_t Z = \partial_x^2 Z + \lambda Z \xi . \quad (\star)$$

Idea: Take this as **definition** of solution, where (\star) is interpreted in the **Itô** sense. Work by Bertini-Giacomin shows that this is the **physically relevant** solution.

Write $h = \mathcal{S}_{CH}(h_0, \omega)$, taking values in $\mathcal{C}(\mathbf{R}_+, \mathcal{C})$.

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Properties of Cole-Hopf

Mollify W , so $W_{\varepsilon,k} = \varphi(\varepsilon k)W_k$ for cutoff φ , and set

$$dZ_\varepsilon = \partial_x^2 Z_\varepsilon dt + \lambda Z_\varepsilon dW_\varepsilon, \quad h_\varepsilon = \lambda^{-1} \log Z_\varepsilon.$$

Then h_ε solves

$$\partial_t h_\varepsilon = \partial_x^2 h_\varepsilon + \lambda((\partial_x h_\varepsilon)^2 - C_\varepsilon) + \xi_\varepsilon, \quad C_\varepsilon \approx \frac{1}{\varepsilon} \int \varphi^2.$$

Problems with this notion of solution:

1. Not satisfactory at the formal level.
2. **Lack of robustness:** no good approximation theory for other modifications (hyperviscosity, time-smoothing, etc).
3. Properties of solutions do **not always transform well** (regularity of difference for example).

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Some attempts

1. Consider product as **Wick product** $\partial_x h \diamond \partial_x h$ (Øksendal & Al 1995): **wrong notion** of solution ($\neq \mathcal{S}_{CH}$). Also **wrong scaling properties** (Chan 2000).
2. Formulate as **martingale problem** (Assing 2002): **no well-posedness**, “generator” not shown to be closable.
3. Apply “standard” **renormalisation techniques** inspired by QFT (Da Prato, Debussche, Tubaro 2007): only works for a **regularised equation**.
4. Define nonlinearity on some **distributional space** (Gonçalves, Jara 2010, Assing 2011): **no uniqueness**. No characterisation of class of distributions for which formulation even makes sense.

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A robustness result

Theorem (H. 2011): For $\alpha > 0$ arbitrary one can build the following objects:

$$\begin{array}{ccc} \mathcal{X} \times \mathcal{C}^\alpha & \xrightarrow{\mathcal{S}_R} & \mathcal{C}(\mathbf{R}_+, \mathcal{C}^\alpha) \\ \Psi \uparrow & & \downarrow \\ \Omega \times \mathcal{C}^\alpha & \xrightarrow{\mathcal{S}_{CH}} & \mathcal{C}(\mathbf{R}_+, \mathcal{C}^\alpha) \end{array}$$

where \mathcal{S}_R is **jointly continuous**, but Ψ is only measurable. (Slight cheat: solutions are only local in general.)

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A deterministic result

Consider solutions h_ε to

$$\partial_t h_\varepsilon = \partial_x^2 h_\varepsilon + (\partial_x h_\varepsilon)^2 + \varepsilon^{-3/2} g(\varepsilon^{-1}x - \varepsilon^{-2}t) - K_\varepsilon ,$$

for centred **periodic** g and **suitable large constants** K_ε . Then, one can compute K such that $h_\varepsilon \rightarrow h$ solving

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Proof: Just show that $\Psi(g_\varepsilon)$ converges to a limit in \mathcal{X} ...

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Ideas of technique

Idea: Perform **Wild expansion** of solution: define

$$\partial_t Y_\varepsilon^\bullet = \partial_x^2 Y_\varepsilon^\bullet + \xi_\varepsilon .$$

For any **binary tree** $\tau = [\tau_1, \tau_2]$, define Y_ε^τ recursively by

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$$h_\varepsilon(t) = \sum_{\tau} \lambda^{|\tau|-1} Y_\varepsilon^\tau(t) ,$$

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A convergence result

Theorem: For every τ , there is a choice of α_τ and C_ε^τ such that

$$Y_\varepsilon^\tau \rightarrow Y^\tau \quad (\text{Independent of } \varphi.)$$

in probability in $\mathcal{C}(\mathbf{R}, \mathcal{C}^\alpha) \cap \mathcal{C}^\beta(\mathbf{R}, \mathcal{C})$ for $\alpha < \alpha_\tau$ and $\beta < \frac{1}{2}$.

Optimal choice: $\alpha_\bullet = \frac{1}{2}$, $\alpha_V = 1$, $\alpha_\tau = (\alpha_{\tau_1} \wedge \alpha_{\tau_2}) + 1$.

$$C_\varepsilon^V = C_\varepsilon \sim \frac{1}{\varepsilon} \int_{\mathbf{R}} \varphi^2(x) dx ,$$

$$C_\varepsilon^{\mathbb{W}} = \frac{4\pi}{\sqrt{3}} |\log \varepsilon| - C(\varphi) ,$$

$$C_\varepsilon^{\mathbb{W}} = -\frac{1}{4} C_\varepsilon^{\mathbb{W}} .$$

Truncated expansion

Idea: Write h_ε as

$$h_\varepsilon = \sum_{\tau \in \mathcal{T}} \lambda^{|\tau|-1} Y_\varepsilon^\tau + u_\varepsilon ,$$

for a **finite** set \mathcal{T} , derive an equation for u_ε , and pass to limit.

Minimal working choice: $\mathcal{T} = \{\bullet, \mathfrak{v}, \mathfrak{V}, \mathfrak{W}, \mathfrak{V}\}$. One obtains

$$\partial_t u_\varepsilon = \partial_x^2 u_\varepsilon + 2\lambda \partial_x u_\varepsilon \partial_x Y_\varepsilon^\bullet + \text{"l.o.t."} .$$

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How to solve that equation?

Writing $v = \partial_x u$, recall we want to solve

$$\partial_t v = \partial_x^2 v + 2\lambda \partial_x (v \partial_x Y^\bullet).$$

If v were **constant** on the right hand side, then one would expect v to “look locally like” $2\lambda v \Phi$, where

$$\partial_t \Phi = \partial_x^2 \Phi + \partial_x^2 Y^\bullet.$$

Idea: Set up fixed point argument in space of functions that “look like Φ ” and use the fact that one can define $\Phi \partial_x Y^\bullet$ “by hand”.

Resulting space is a **non-linear** algebraic variety embedded in a larger Banach space. Uses **controlled rough paths** à la Gubinelli-Lyons.

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Take-away message: Nonlinear spaces are **required** to solve some rough equations **pathwise**.

Some open problems:

- Extension to $x \in \mathbb{R}$?
- Convergence of microscopic models (for example lattice KPZ) to KPZ. See work with J. Maas and H. Weber.
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- Rough equations in higher dimensions?

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